

Aggregation–Diffusion Equations: concentration vs simplification

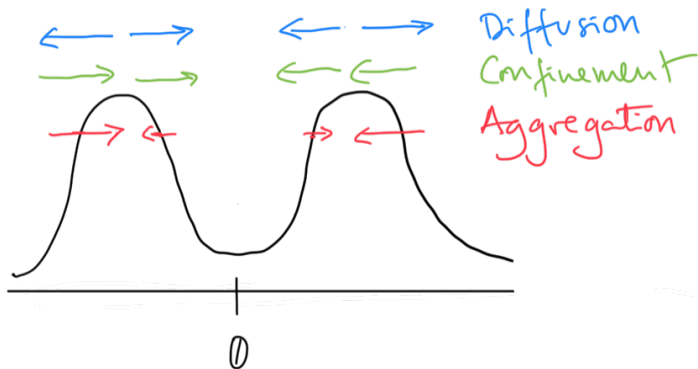
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Febrero 2023

Aggregation and Diffusion



A model of diffusion, confinement and aggregation

The aim of this talk is to explain the modeling and theory behind the following model for aggregation-diffusion phenomena:

$$\frac{\partial \rho}{\partial t} = \operatorname{div} \left(\rho \nabla \left(\underbrace{U'(\rho)}_{\text{Diffusion}} + \underbrace{V}_{\text{Confinement}} + \underbrace{W * \rho}_{\text{Aggregation}} \right) \right) \quad (\text{ADE})$$

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We will discuss the range of power-type aggregation and diffusion

$$U'(\rho) = \frac{m}{m-1} \rho^{m-1}, \quad V(x) = \frac{|x|^\alpha}{\alpha}, \quad \text{and} \quad W(x) = \frac{|x|^\lambda}{\lambda}.$$

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If V, W are bounded below, we can always assume $V, W \geq 0$.

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Classical results of asymptotics

Calculus of Variations approach

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Minimisation for ADE and asymptotic concentration

An example of asymptotic simplification

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Notice $\Delta \varphi(\rho) = \operatorname{div}(\varphi'(\rho) \nabla \rho)$ so $U''(\rho) = \frac{\varphi'(\rho)}{\rho}$.

Particle systems

Consider N with positions X_i of equal masses $1/N$

¹Assume $\nabla W(0) = 0$

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Consider N with positions X_i of equal masses $1/N$ and the attracting/repulsive system¹

$$\frac{dX_i}{dt} = - \sum_{\substack{j=1 \\ j \neq i}}^N \underbrace{\frac{1}{N} \nabla W(X_i - X_j)}_{\text{Aggregation}} - \underbrace{\frac{1}{N} \nabla V(X_i)}_{\text{Confinement}}, \quad i = 1, \dots, N$$

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Diffusion can be added to the particle system by introducing noise

[► Details](#)

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The Aggregation-Diffusion Equation

Joining the many particle approximation with the Porous Medium diffusion:

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Keller-Segel ($d = 2$)	$\rho \log \rho$	0	$-\frac{1}{2\pi} \log x $
Swarming / Herding	0	0	$\frac{1}{a} x ^a - \frac{1}{b} x ^b$

Conservation?

In conservation laws, we expect $\int_{\mathbb{R}^d} \rho(t) = \int_{\mathbb{R}^d} \rho_0$
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$$\frac{d}{dt} \int_{\mathbb{R}^d} \rho \, dx = \frac{d}{dt} \lim_{R \rightarrow \infty} \int_{B_R} \rho \, dx = \lim_{R \rightarrow \infty} \int_{\partial B_R} j \frac{x}{|x|} \, dS \stackrel{?}{=} 0.$$

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Sometimes mass is not conserved, and we will give an example later.

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It admits a solution

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- ▶ Any solution satisfies

$$\|\rho(t, \cdot) - K(t, \cdot)\|_{L^1} \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

This is known as **asymptotic simplification**.

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- [Brezis and Friedman 1983] proved that δ_0 does not diffuse: if we take a sequence $\rho_j(0^+, \cdot) \rightarrow \delta_0$, the associated solution $\rho_j(t) \rightarrow \delta_0$.

Asymptotic profiles

Notice that the heat kernel is **self-similar**

$$K(t, x) = (2t)^{-\frac{d}{2}} G\left(-\frac{x}{\sqrt{2t}}\right), \quad G(y) = \frac{1}{(2\pi)^{\frac{d}{2}}} \exp\left(-\frac{|y|^2}{2}\right)$$

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Furthermore, G is an **asymptotic profile** for the equation:

$$\|u(t, \cdot) - G\|_{L^1} \rightarrow 0, \quad \text{as } t \rightarrow \infty.$$

Asymptotic profiles

Notice that the heat kernel is **self-similar**

$$K(t, x) = (2t)^{-\frac{d}{2}} G\left(-\frac{x}{\sqrt{2t}}\right), \quad G(y) = \frac{1}{(2\pi)^{\frac{d}{2}}} \exp\left(-\frac{|y|^2}{2}\right)$$

So send K to G : $\tau = \log \sqrt{2t+1}$, $y = \frac{x}{\sqrt{2t+1}}$ and

$$u(\tau, y) = e^{d\tau} \rho(t, x)$$

Applying the change of variable to the heat equation we recover the **Fokker-Planck equation**

$$\begin{aligned} \partial_\tau u &= \Delta_y u + \operatorname{div}(yu) \\ &= \operatorname{div}(u \nabla_y (\log u + \frac{|y|^2}{2})). \end{aligned}$$

Clearly, G is a stationary solution.

Furthermore, G is an **asymptotic profile** for the equation:

$$\|u(t, \cdot) - G\|_{L^1} \rightarrow 0, \quad \text{as } t \rightarrow \infty.$$

A similar approach works for the Porous Medium Equation, where the profile is B .

Finite time blow-up

The Keller-Segel model

The Keller-Segel proposed a model of cell migration by chemotaxis given by

$$\begin{cases} \partial_t \rho = \Delta \rho + \operatorname{div}(\rho \nabla v), \\ -\Delta v = \rho. \end{cases} \quad M = \int_{\mathbb{R}^d} \rho_0(x) \, dx$$

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There exists $M^* > 0$ such that

- ▶ If $M < M^*$ solutions are global-in-time.
- ▶ If $M > M^*$ there is finite-time blow-up. And $\rho(T^*) = M\delta_0$.

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When $M > M^*$ there exist $\|\rho_0\|_{L^1} = M$ such that $\mathcal{F}[\rho_0] < 0$.

Main question

For

$$\frac{\partial \rho}{\partial t} = \operatorname{div} \left(\rho \nabla (U'(\rho) + V + W * \rho) \right) \quad (\text{ADE})$$

can we classify characterise ρ_∞ such that

$$\rho(t) \rightarrow \rho_\infty$$

in terms of general U, V, W ?

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The Heat Equation as a gradient flow in L^2

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Remark

We can rewrite the Heat Equation

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In general, the $\nabla_{L^2} \mathcal{F}$ is given by the Euler-Lagrange equations

Gradient flows in Wasserstein space

Our equations are “nice” in 2-Wasserstein space (\mathcal{P}_2).

For $\mathcal{F} : L^1 \cap \mathcal{P}_2(\mathbb{R}^d) \rightarrow \mathbb{R}$ formally speaking

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If $W(x) = W(-x)$, we can formally rewrite the Aggregation-Diffusion problem as²

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²Due to the convolution, \mathcal{F} is non-local and $\mathcal{F}[\rho] \neq \int_{\mathbb{R}^d} F(x, \rho(x)) \, dx$.
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Formally, $\frac{d}{dt} \mathcal{F}[\rho(t)] = - \int_{\mathbb{R}^d} \rho \left| \nabla \frac{\delta \mathcal{F}}{\delta \rho} [\rho] \right|^2.$

This is called *energy dissipation* estimate.

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The precise definition of solution is the notion of *curves of maximal slope*.

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Gradient-flow structure and minimisation

The extension of convexity in \mathbb{R}^d for \mathcal{P}_2 is **displacement convexity** (see [McCann 1997])

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Due to the energy dissipation a minimiser ρ_∞ should satisfy

$$\rho_\infty \nabla \frac{\delta \mathcal{F}}{\delta \rho} [\rho_\infty] = 0.$$

Either $\rho_\infty = 0$ (as in PME), or $\frac{\delta \mathcal{F}}{\delta \rho} [\rho_\infty] = C$ (over open sets).

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This gives the intuition that for $M > 8\pi$ then δ_0 (i.e. $\lambda \rightarrow \infty$) is energy beneficial.

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We go back to

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$$\rho(h) = \left(\frac{1-m}{m} (V + h) \right)^{\frac{-1}{1-m}}$$

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We are able to prove, in some cases, that they are global attractors of (ADE): talk of A. Fernández-Jiménez.

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Therefore, in this range we always expect diffusion.

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Then (\star) .

Notice that this hypothesis work for $W(x) \sim |x|^{-\varepsilon}$ for any $\varepsilon > 0$, but not for the critical case $W(x) \sim \log |x|$.

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- ▶ Thus, using the control of $\int \rho \log \rho$ we prove that

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as $t \rightarrow \infty$.



Questions, comments, remarks, ...

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